

Nucleon structure from $N_f = 2 + 1$ DWF simulations

Takeshi YAMAZAKI



T. Blum, H-W. Lin, S. Ohta, K. Orginos, and S. Sasaki
for RBC-UKQCD Collaboration

Domain Wall Fermions at Ten Years @ BNL
March 15-17, 2007

Outline

1. Introduction
2. Simulation parameters
3. Preliminary results
 - g_A/g_V
 - Moments of quark distributions
 - Form factors
4. Summary

1. Introduction

Motivation :

understand nucleon structure from first principle

We calculate matrix elements related to nucleon structure on $N_f = 2+1$ DWF configuration.

- g_A/g_V
Well determined experimentally: $g_A/g_V = 1.2673(35)$
- Moments of quark distributions
Deep inelastic scattering; structure functions
 $\langle x \rangle_q \rightarrow$ Unpolarized: $F_1(x, Q^2), F_2(x, Q^2)$
 $\langle x \rangle_{\Delta q} \rightarrow$ Polarized: $g_1(x, Q^2), g_2(x, Q^2)$
- Form factors
Elastic scattering

$$\begin{aligned} F_1(q^2) &= \frac{1}{(1 + q^2/M_V^2)^2}, & \langle r_{ch}^2 \rangle &= 12/M_V^2, \\ G_A(q^2) &= \frac{g_A}{(1 + q^2/M_A^2)^2}, & \langle r_{ax}^2 \rangle &= 12/M_A^2 \end{aligned}$$

RBC-UKQCD generated $N_f = 2 + 1$ dynamical configuration.

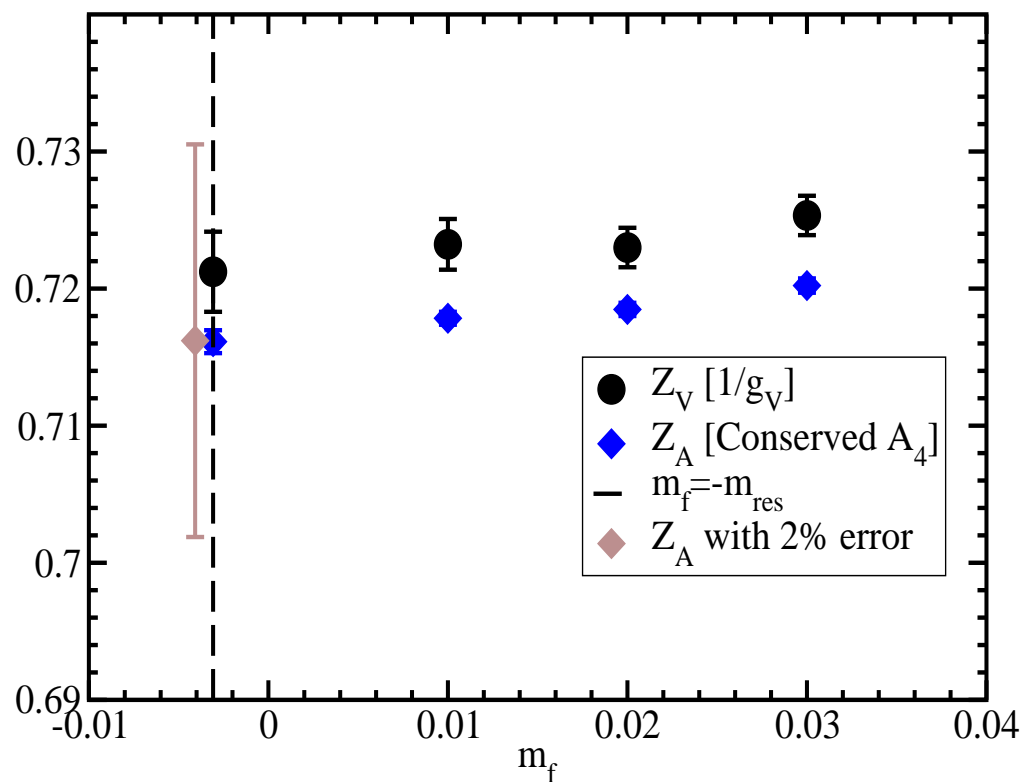
1. u, d quark mass is as low as $m_\pi = 310$ MeV.

Investigation of nucleon structure in chiral regime

2. Chiral symmetry on lattice

3. Physical volume 3 fm

Chiral symmetry on lattice



If chiral symmetry is exact,
 $Z_V = Z_A$ at chiral limit.

Z_V is determined by
Nucleon form factor

$$Z_V g_V^{\text{lat}} = F_1(0) = 1$$

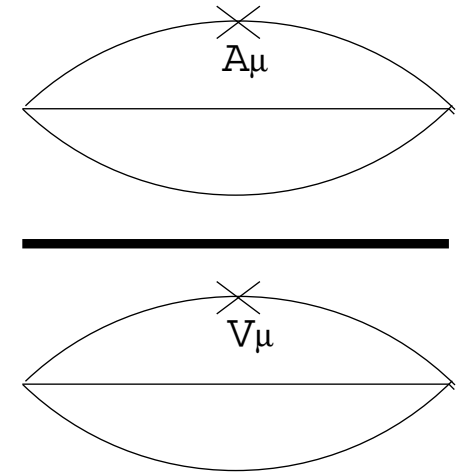
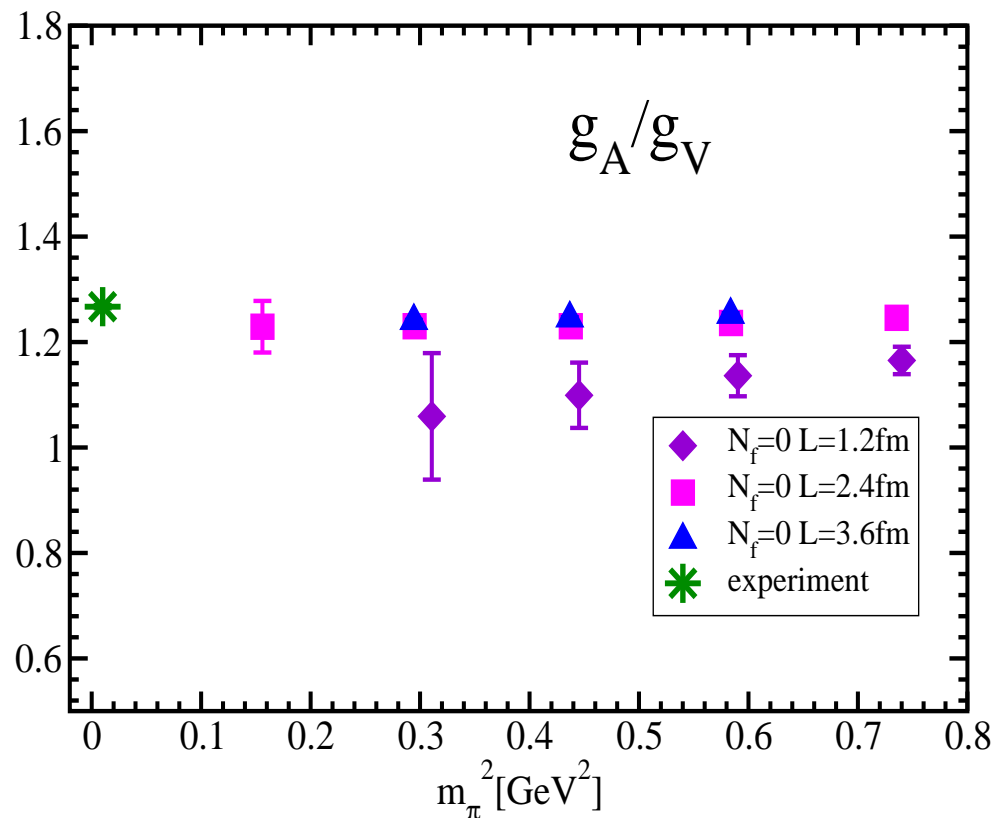
Z_A is determined by
Conserved axial-vector current \mathcal{A}_0

Z_A is consistent with Z_V within 2% at chiral limit.

Finite volume effect of nucleon matrix element

g_A/g_V is a simple, basic physical quantity of nucleon structure.

It is easy to calculate with DWF due to $Z_V/Z_A \approx 1$.



Large finite volume effect is seen in heavy m_π region.

g_A/g_V at 2.4 fm almost agrees with one at 3.6 fm.

$L \approx 2.5$ fm is enough for nucleon calculation.

RBC-UKQCD generated $N_f = 2 + 1$ dynamical configuration.

1. u, d quark mass is as low as $m_\pi = 310$ MeV.

Investigation of nucleon structure in chiral regime

2. Chiral symmetry on lattice

$Z_V = Z_A$ is satisfied within a few %.

3. Physical volume 3 fm

Volume is large enough for nucleon calculation, based on $N_f = 0$ calculation.

3. Simulation parameters

- $N_f = 2 + 1$ Iwasaki gauge + Domain Wall fermion actions
- $\beta = 2.13$ $a^{-1} = 1.62$ GeV $M_5 = 1.8$ $m_{\text{res}} = 0.003$
- Lattice size $24^3 \times 64 \times 16$ ($La \approx 3$ fm)
- $m_s = 0.04$ fixed (close to m_s^{phys})
- quark masses and confs.

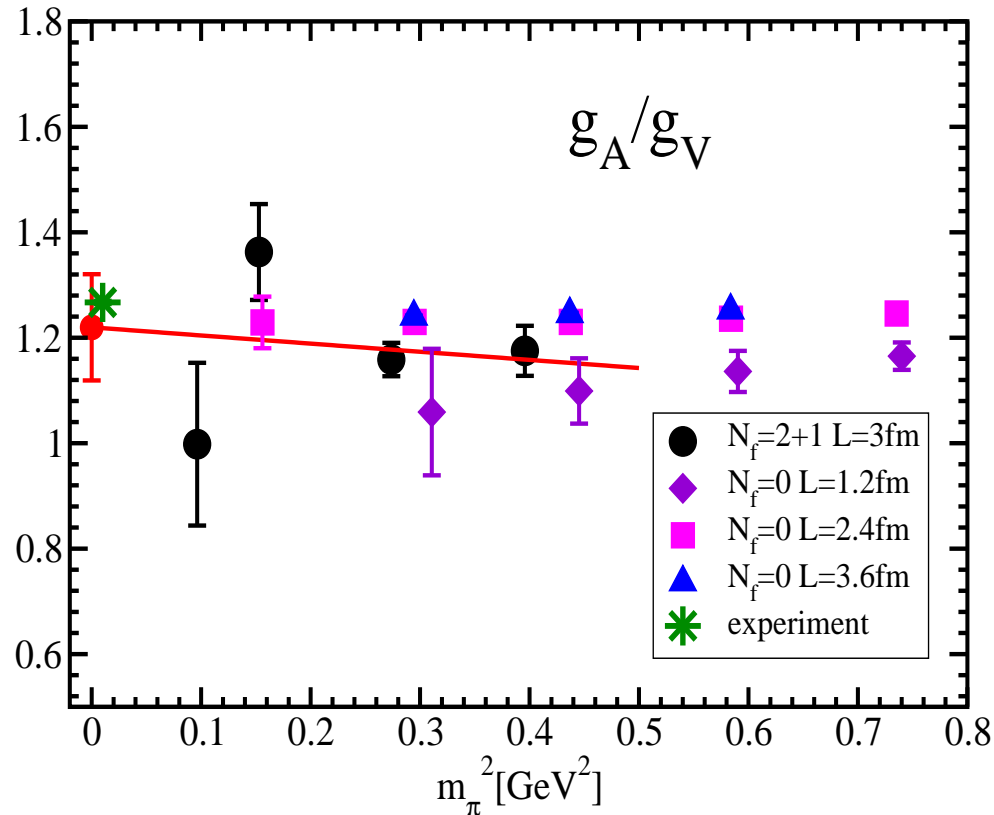
m_f	m_π [MeV]	# of confs. $\times N_{\text{meas}}$
0.005	310	52×4
0.01	390	119×4
0.02	520	49×4
0.03	690	53×4

Results at two lighter masses do not have good accuracy, so that
all results are preliminary.

- We focus only on iso-vector quantities. (no disconnected diagram)

4. Preliminary results

4.1. g_A/g_V



Preliminary result

$$g_A/g_V = \begin{cases} 1.22(10) & (\text{lat.}) \\ 1.267(4) & (\text{exp.}) \end{cases}$$

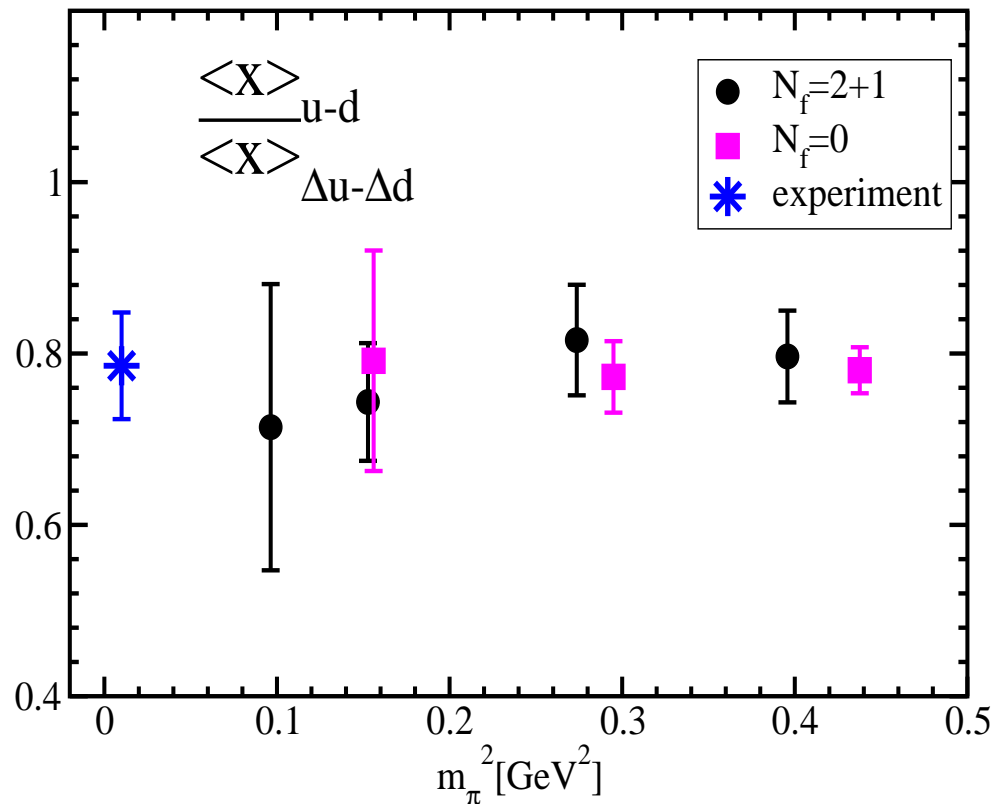
We will confirm the result is reliable by improving statistics.

m_π is lighter than m_π in $N_f = 0$ case.

Results at two lighter m_π has larger error and fluctuation.

There seems to be no dynamical effect.

4.2. Moments of quark distributions



Unpolarized : $\langle x \rangle_{u-d}$

Polarized : $\langle x \rangle_{\Delta u - \Delta d}$

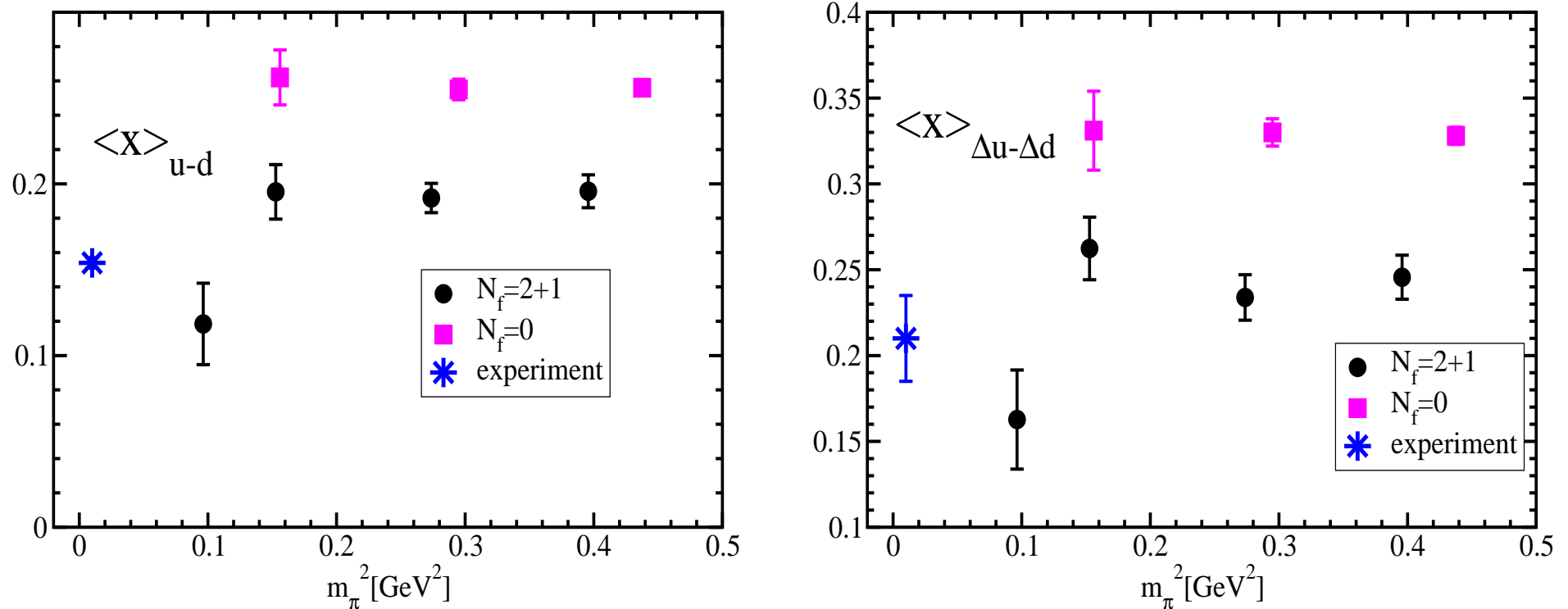
If chiral symmetry is exact, renormalizations of $\langle x \rangle_{u-d}$ and $\langle x \rangle_{\Delta u - \Delta d}$ are same.

Result is consistent with experiment as well as in $N_f = 0$ case.

However, ...

4.2. Moments of quark distributions (cont'd)

Each component in $N_f = 0$ is independent of m_π .



$\langle x \rangle_{u-d}$ and $\langle x \rangle_{\Delta u - \Delta d}$ are closer to experiment, and have some m_π dependence.

Perturbative $Z(2\text{GeV}) = 0.88(5)$ from PLB641,67

→ We will calculate Z factor by non-perturbative method.

Preliminary result

4.3. Form factors

Iso-vector form factors $F_i^p - F_i^n$

$$\langle N, p | V_\mu(q) | N, p' \rangle = \langle N, p | F_1(q^2) \gamma_\mu + i \sigma_{\mu\nu} q_\nu \frac{F_2(q^2)}{2M_N} | N, p' \rangle$$

$$\langle N, p | A_\mu(q) | N, p' \rangle = \langle N, p | G_A(q^2) i \gamma_5 \gamma_\mu + i \gamma_5 q_\mu G_P(q^2) | N, p' \rangle$$

$$q = p' - p, \quad \left(\frac{Lp'}{2\pi} \right)^2 = 0, 1, 2, 3, 4$$

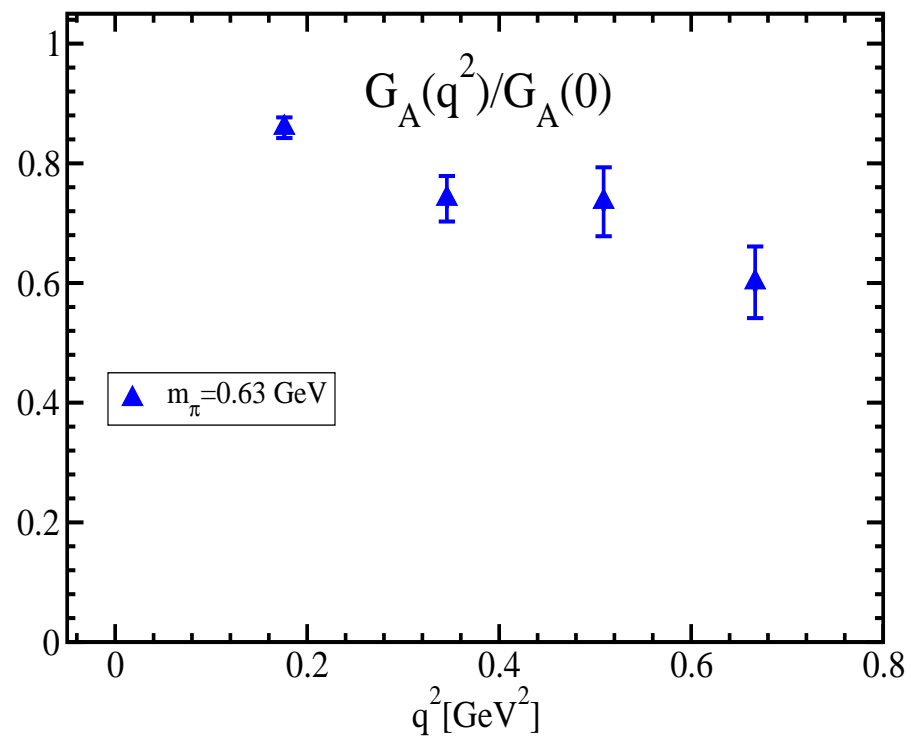
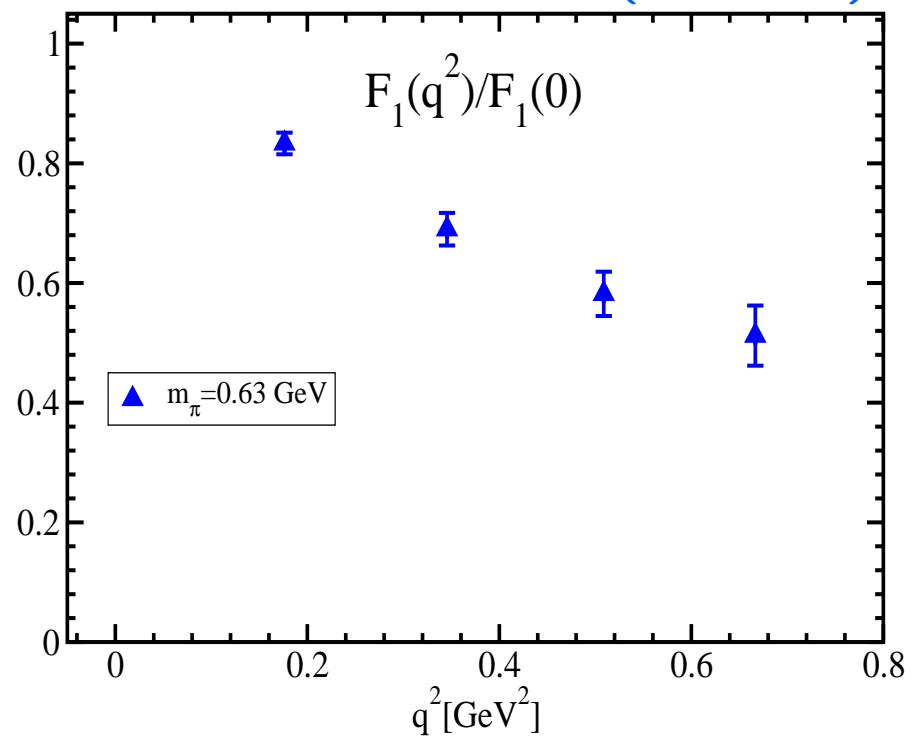
$$F_1(q^2) = \frac{1}{(1 + q^2/M_V^2)^2}, \quad M_V = 0.858(8) \text{ GeV}$$

$$\langle r_{ch}^2 \rangle = 0.636(12) \text{ fm}$$

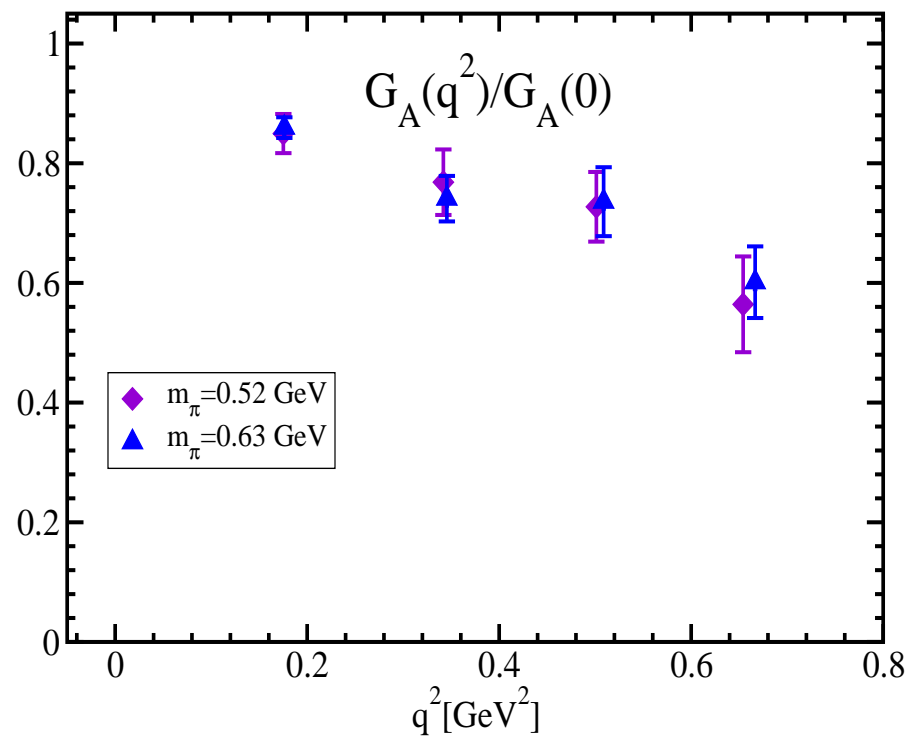
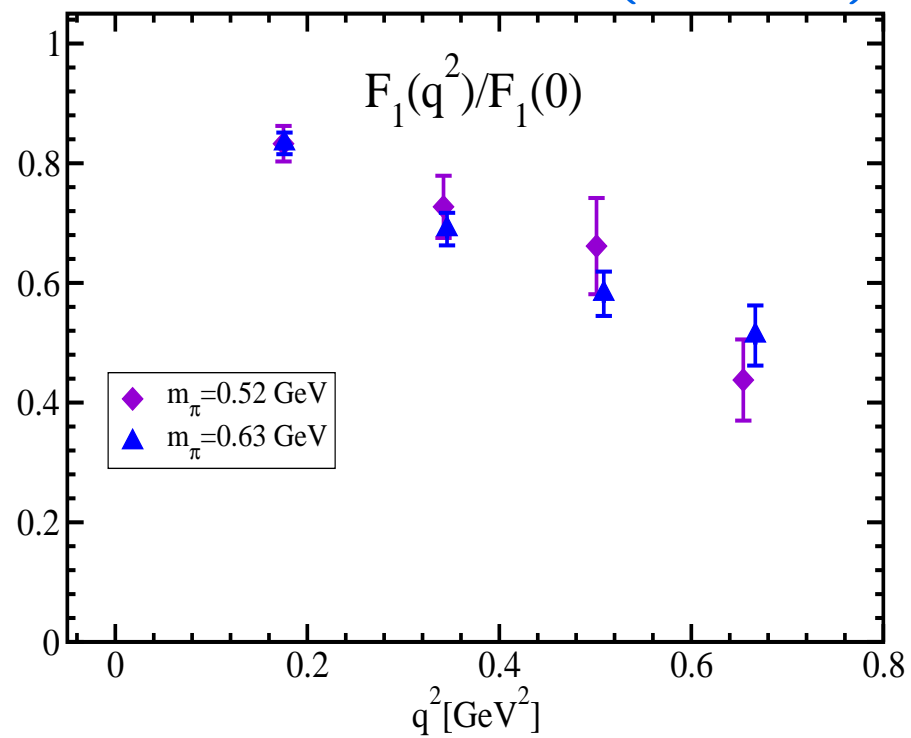
$$G_A(q^2) = \frac{g_A}{(1 + q^2/M_A^2)^2}, \quad M_A = 1.07(2) \text{ GeV}$$

$$\langle r_{ax}^2 \rangle = 0.408(13) \text{ fm}$$

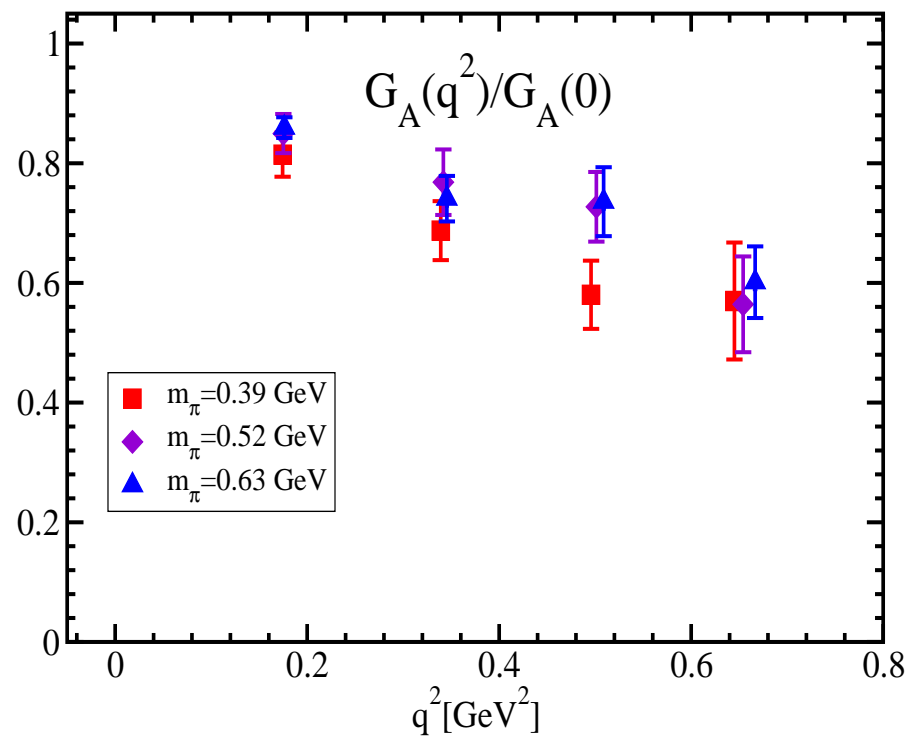
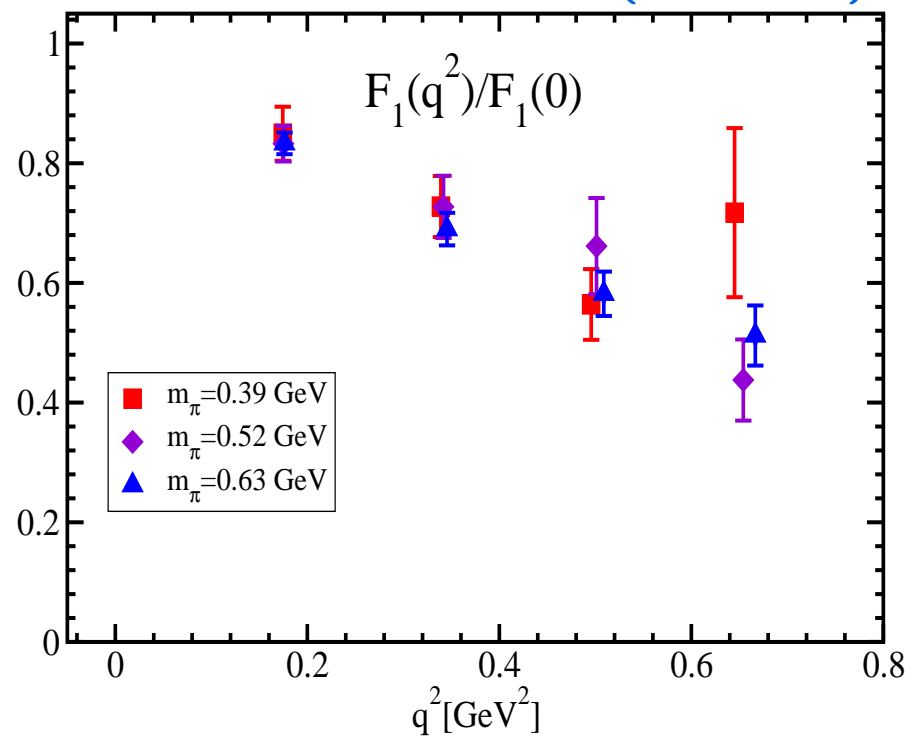
4.3. Form factors (cont'd)



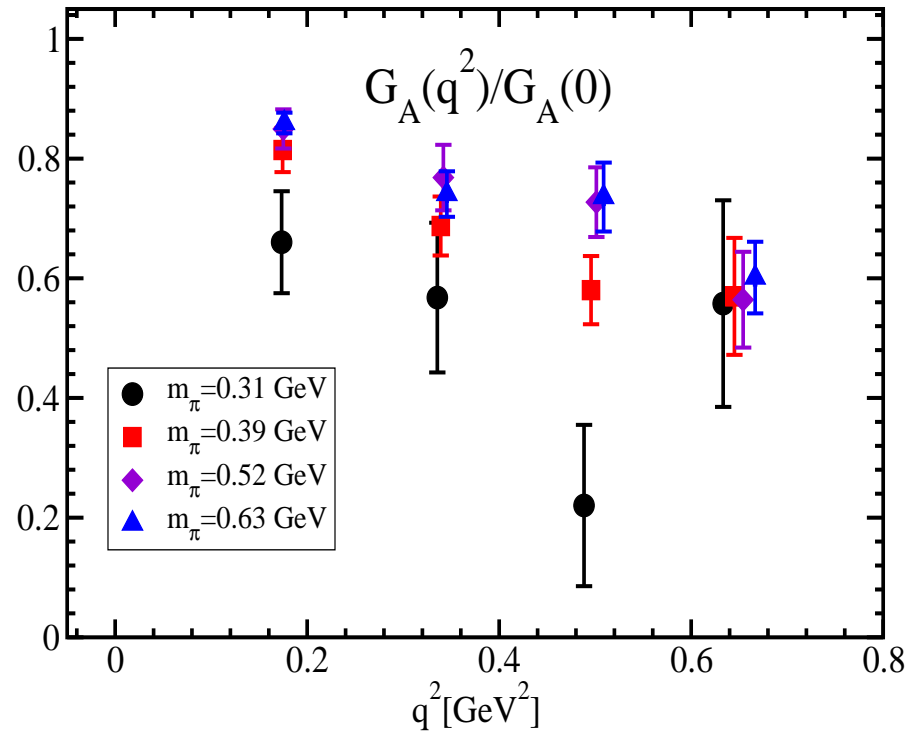
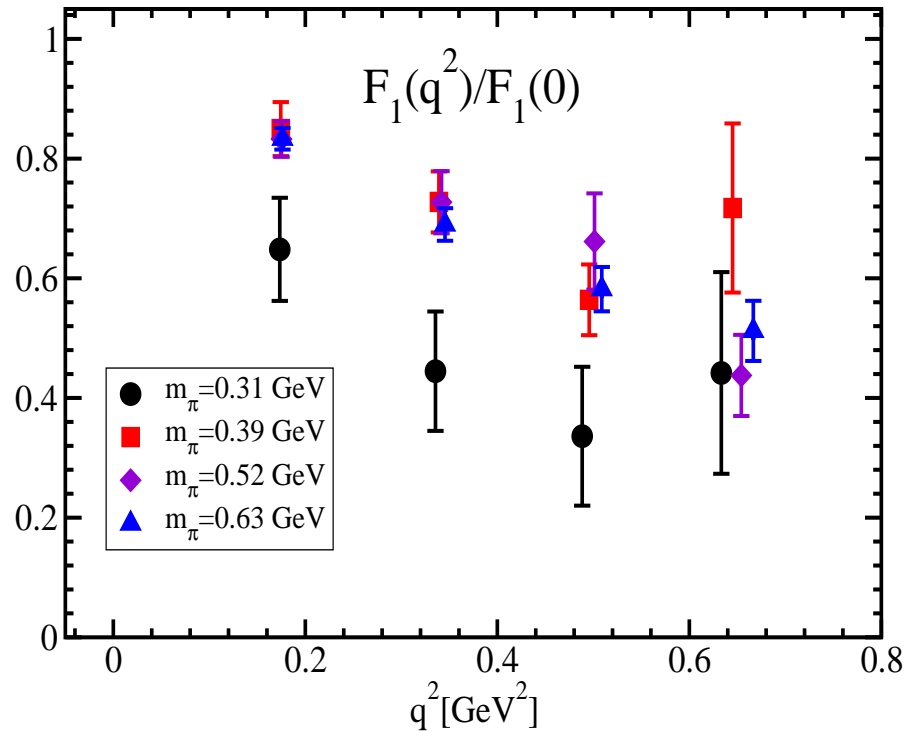
4.3. Form factors (cont'd)



4.3. Form factors (cont'd)



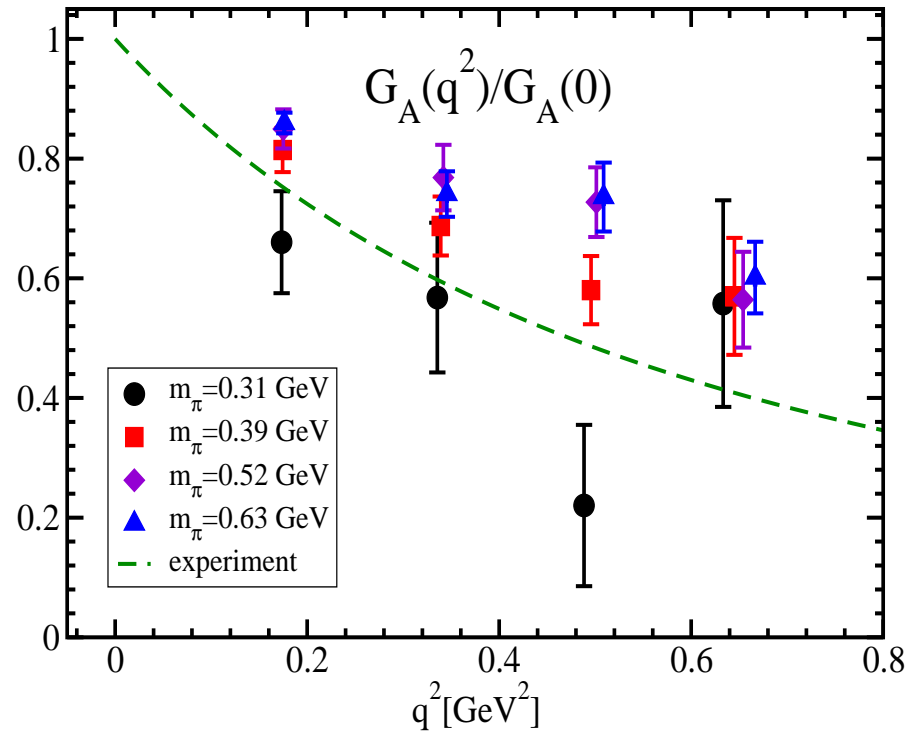
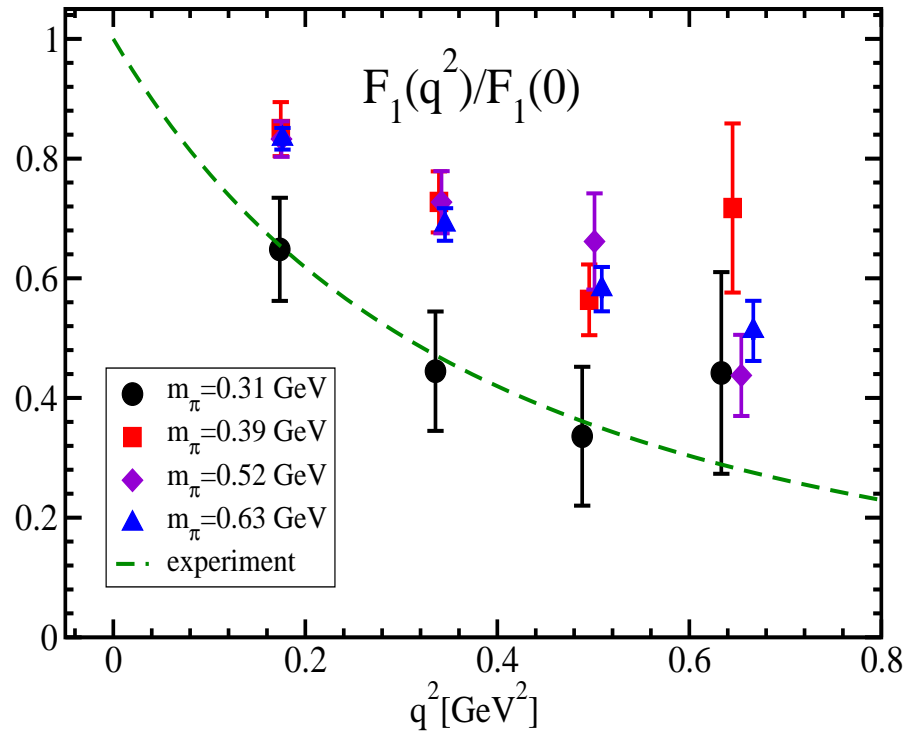
4.3. Form factors (cont'd)



F_1 is almost independent of m_π except for lightest mass.

G_A has m_π dependence.

4.3. Form factors (cont'd)



F_1 is almost independent of m_π except for lightest mass.

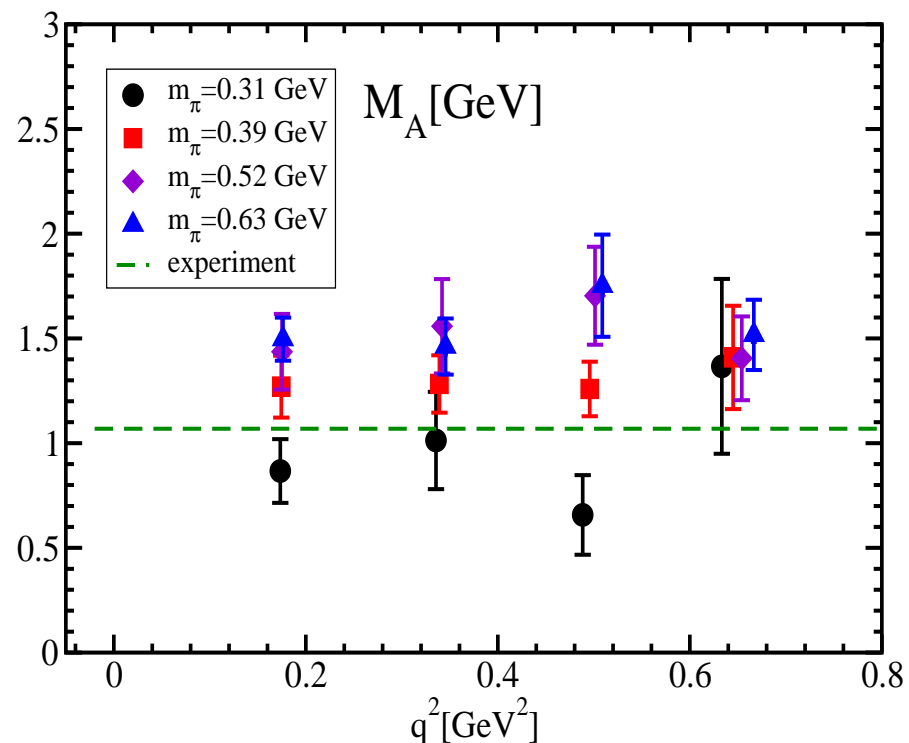
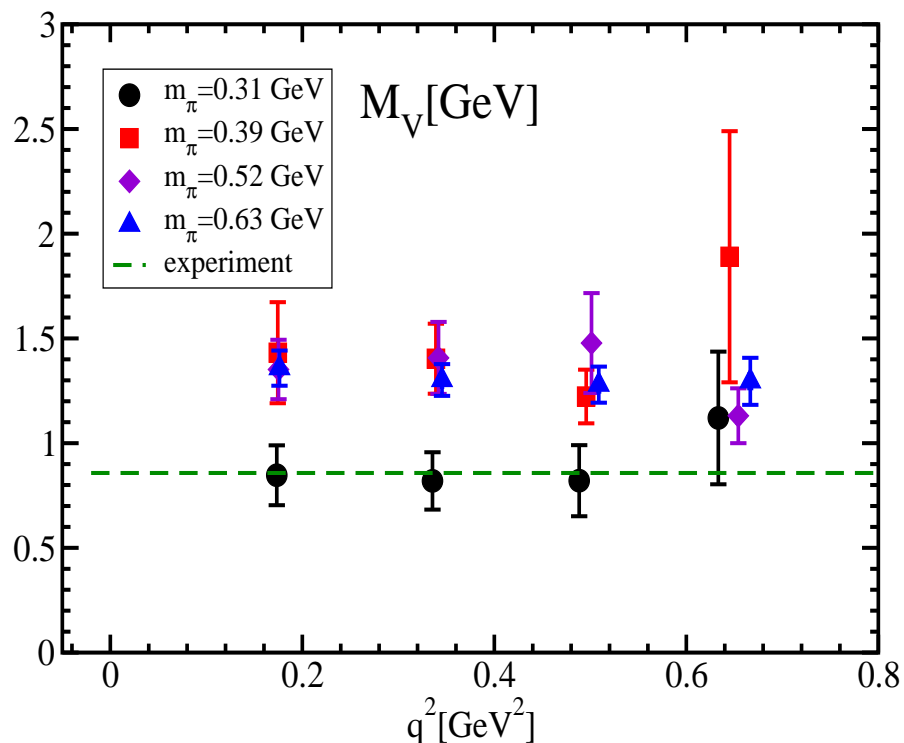
G_A has m_π dependence.

$$F_1(q^2) = \frac{1}{(1 + q^2/M_V^2)^2} \quad G_A(q^2) = \frac{g_A}{(1 + q^2/M_A^2)^2}$$

4.3. Form factors (cont'd)

Effective M_V and M_A Preliminary result

$$M_V = \sqrt{\frac{\sqrt{F_1(q^2)} - 1}{q^2}}, \quad M_A = \sqrt{\frac{\sqrt{G_A(q^2)/g_A} - 1}{q^2}}$$

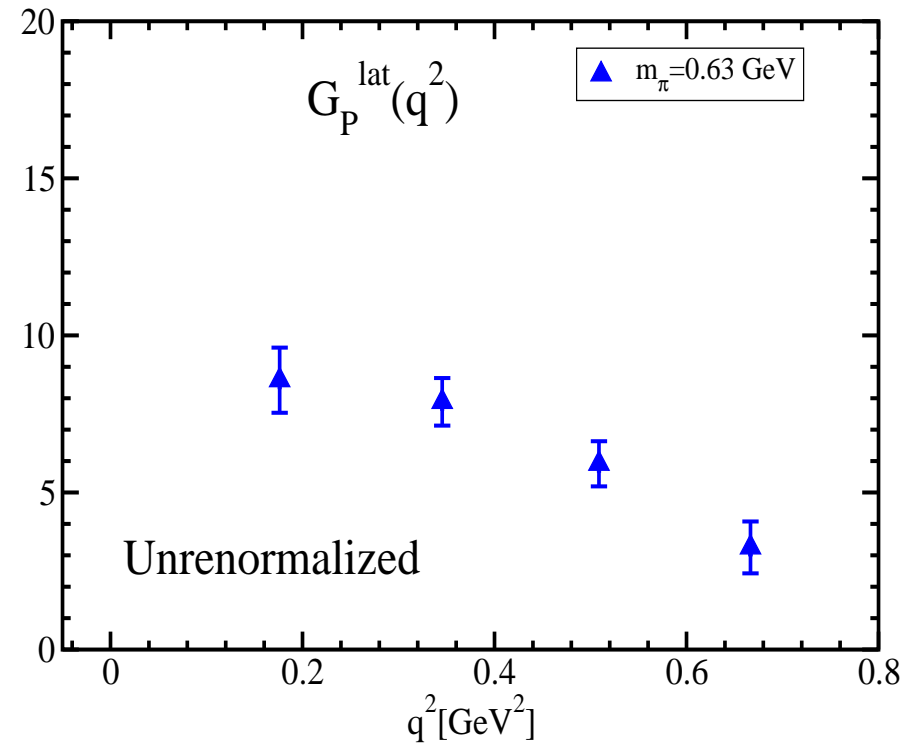
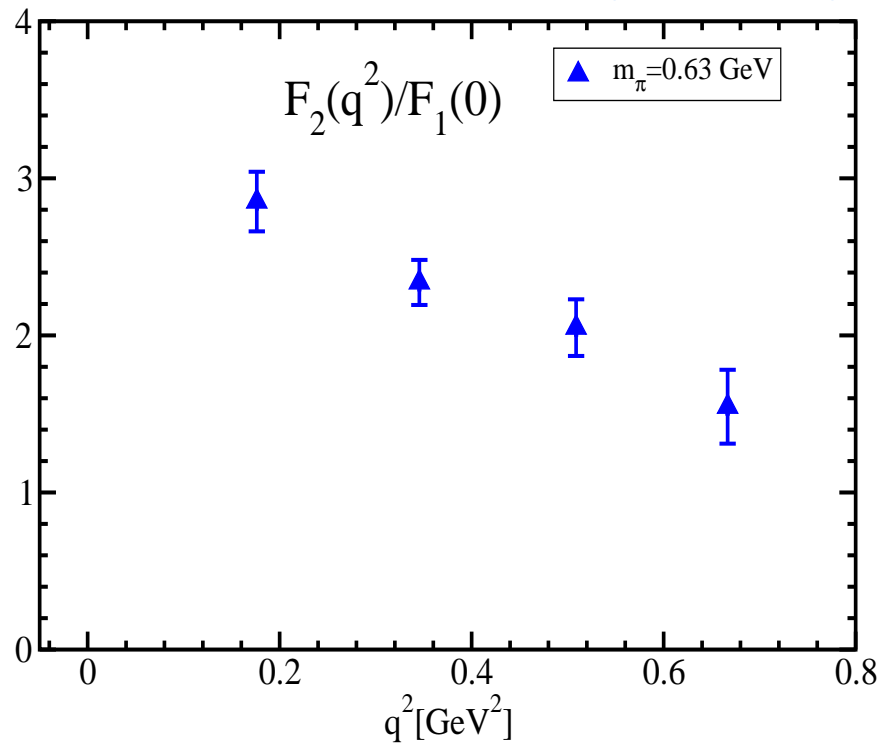


Effective M_V and M_A are reasonably flat.

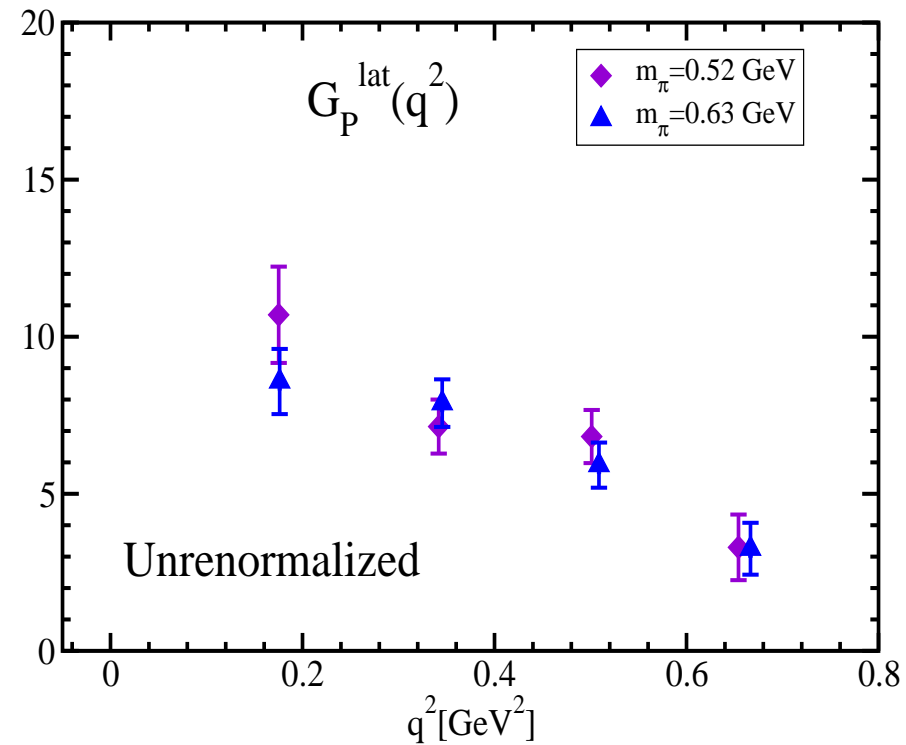
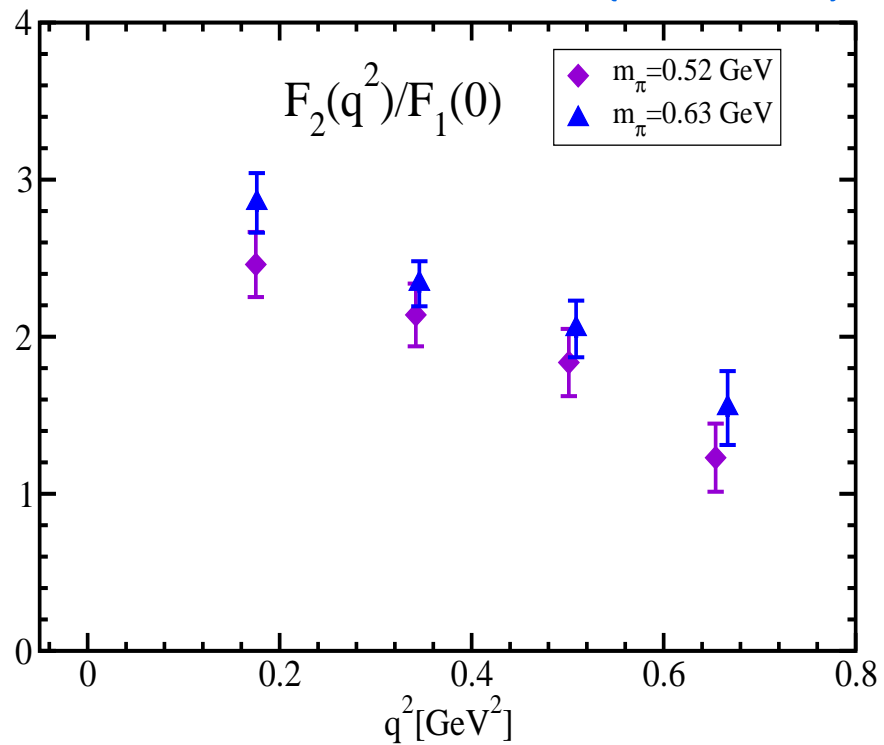
M_V is almost independent of m_π except for lightest mass.

M_A has m_π dependence.

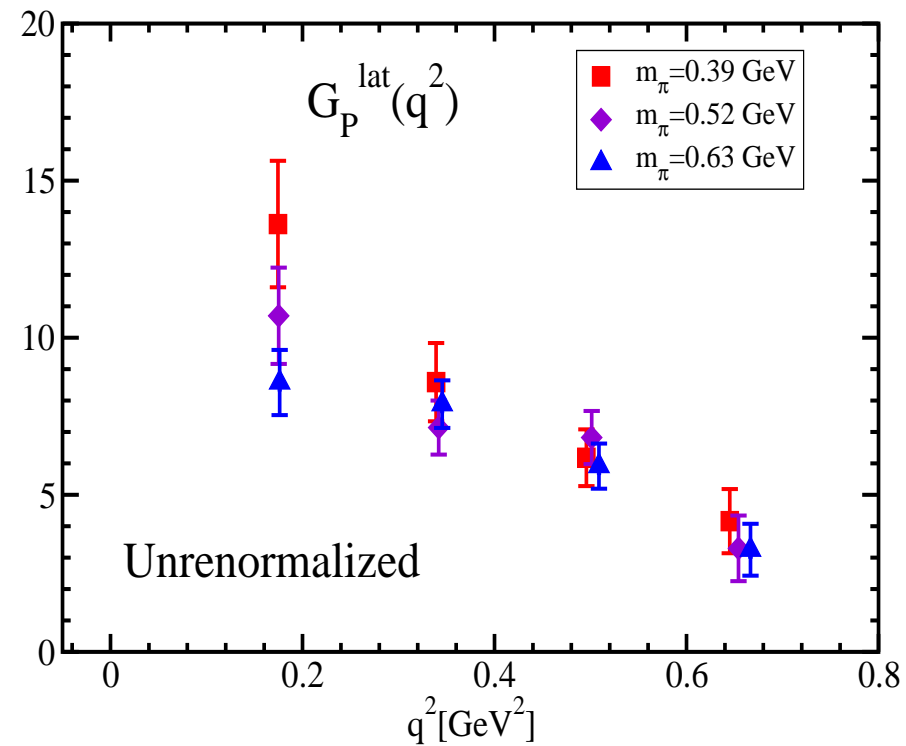
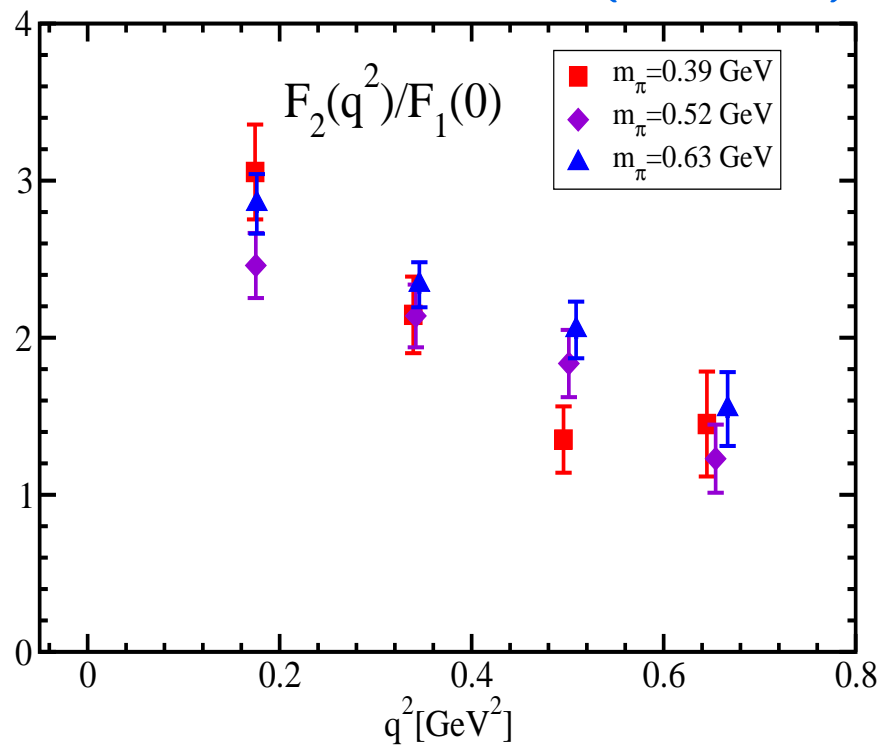
4.3. Form factors (cont'd)



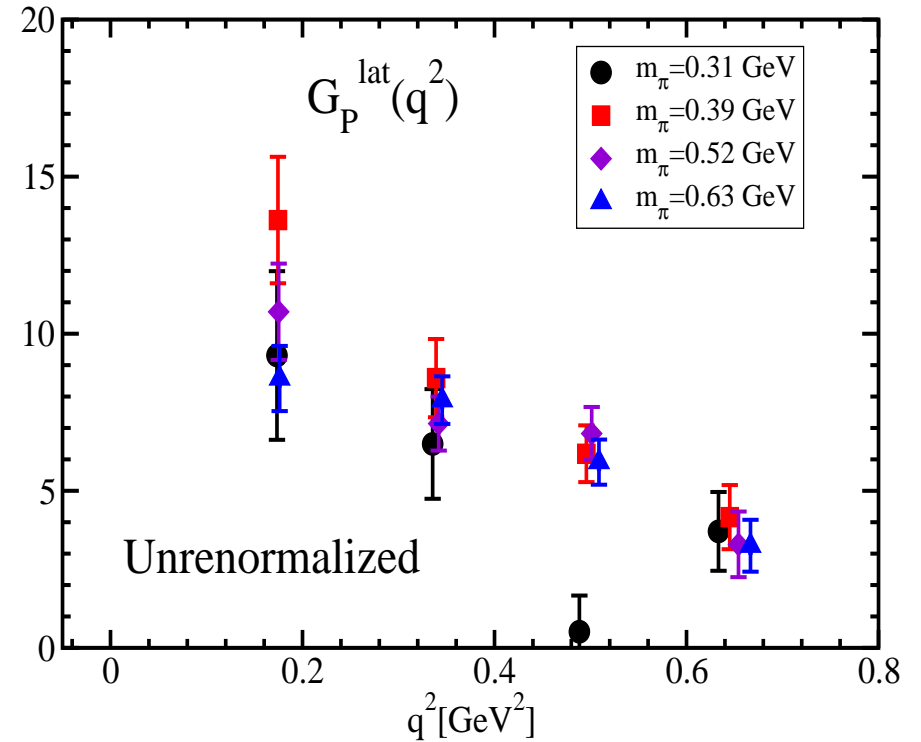
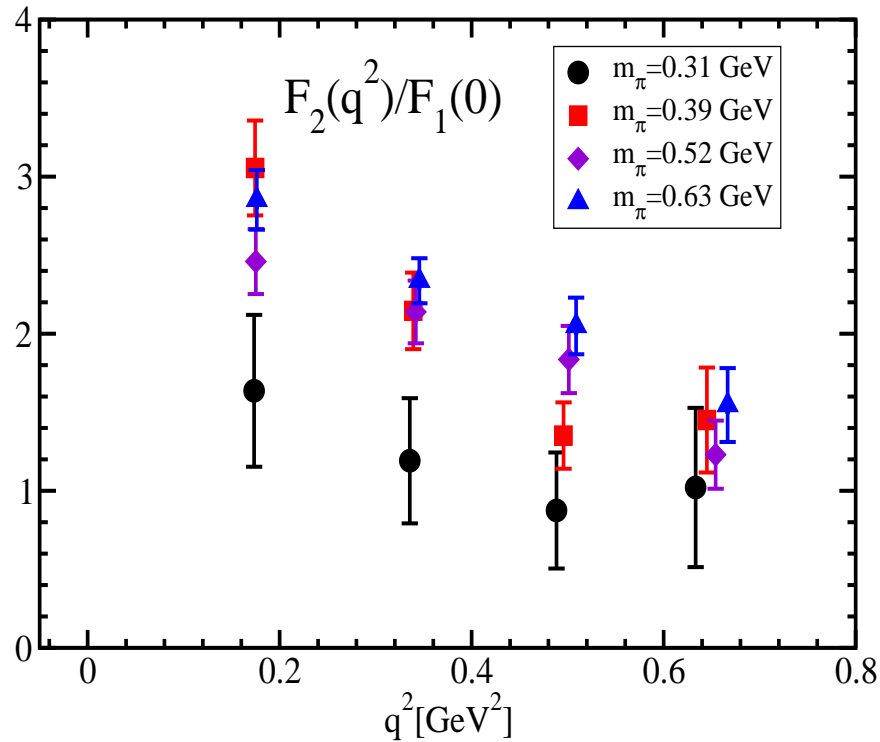
4.3. Form factors (cont'd)



4.3. Form factors (cont'd)



4.3. Form factors (cont'd)

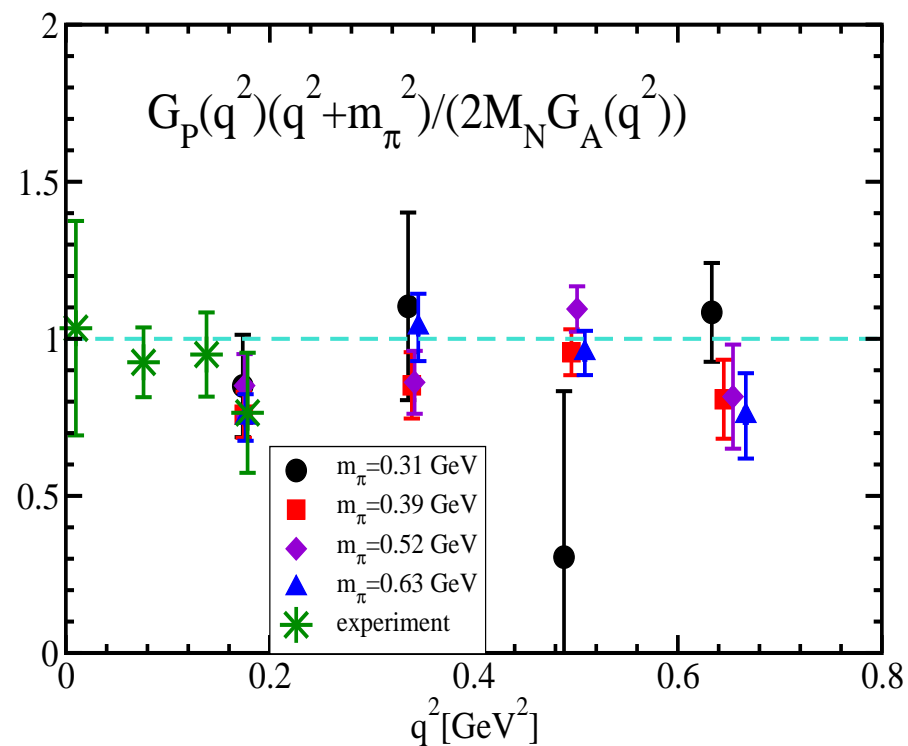
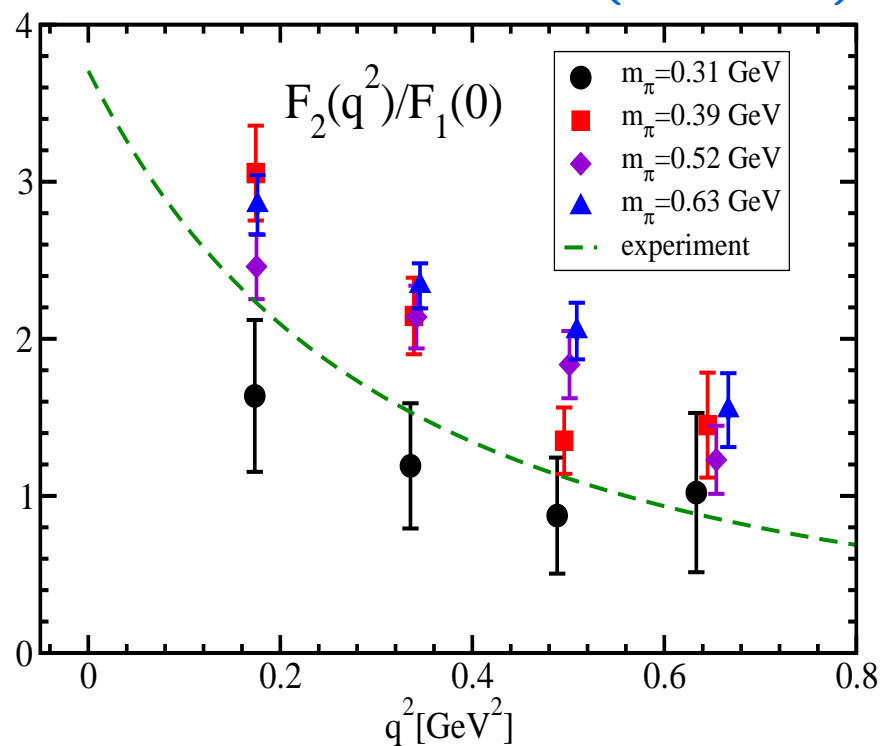


F_2 has m_π dependence, but one at lightest mass is much smaller.

G_P is explained by pion pole in current algebra,

$$G_P(q^2) = \frac{2M_N G_A(q^2)}{q^2 + m_\pi^2}.$$

4.3. Form factors (cont'd)



F_2 at lightest mass is below experiment.

Normalized G_P by pion pole and G_A is almost consistent with current algebra, and close to experiment.

$$G_P(q^2) = \frac{2M_N G_A(q^2)}{q^2 + m_\pi^2}$$

Preliminary result

5. Summary

- We calculated nucleon matrix elements with $N_f = 2 + 1$ dynamical domain wall fermions at light quark masses.
- All results are preliminary.
- We found encouraging and consistent results with experiments.

Future work

- We will improve statistics of $m_f = 0.005$ and $m_f = 0.01$ data.
- Next calculation is on larger lattice size and smaller lattice spacing.